## An operatorial approach to stock markets

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# An operatorial approach to stock markets 

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#### Abstract

We propose and discuss some toy models of stock markets using the same operatorial approach adopted in quantum mechanics. Our models are suggested by the discrete nature of the number of shares and of the cash which are exchanged in a real market, and by the existence of conserved quantities, like the total number of shares or some linear combination of cash and shares. The same framework as the one used in the description of a gas of interacting bosons is adopted.


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## 1. Introduction

A huge literature exists concerning the time behaviour of financial markets, most of which is based on statistical methods, see [1] and references therein. In recent years a somewhat different strategy has also been considered. This strategy takes inspiration from the manybody nature of a stock market, nature which suggests the use of tools naturally related to quantum mechanics and, in particular, to $Q M_{\infty}$, i.e. quantum mechanics for systems with infinite degrees of freedom. Examples of this approach can be found, for instance, in [2] and [3], where the concepts of Hamiltonian, phase transition, symmetry breaking and so on are introduced. However, in none of these papers, and in our knowledge not even in other existing literature, has the analysis of the time evolution of the portfolio of each single trader been undertaken. It should be mentioned, however, that a point of view not very different from the one adopted here is discussed, for instance in [4] and [5].

In this paper we use quantum mechanical ideas to construct some toy models which should mimic a simplified stock market. In all our models, where for simplicity a single kind of share is exchanged, the total number of shares does not change in time. This reminds of what happens in a totally different context, i.e. in a gas of elementary particles which can interact among them but without changing their total number. Also, the price of a single share does not change continuously, since any variation is necessarily an integer multiple of a certain minimal
quantity, the monetary unit, which can be seen, using our quantum mechanical analogy, as a sort of quantum of cash. $Q M_{\infty}$ provides a natural framework in which these features can be taken into account. It also provides some natural tools to discuss the existence of conserved quantities and to find the differential equations of motion which drive the portfolio of each single trader, as we will see.

The paper is organized as follows.
In the next section we discuss a first easy model and we give an interpretation to the quantities used to define the model. This oversimplified model will be useful to fix some general ideas.

In section 3 we improve the model introducing the cash, the price of the share and the supply of the market. We prove that many integrals of motion exist. The equations of motion are solved using a perturbative expansion, well known in $Q M_{\infty}$.

In section 4 we consider a particular version of this model which we completely solve using the so-called mean-field approximation. We also discuss the role of KMS-like states in our framework.

Section 5 contains our conclusions and plans for the future, while in appendix A we give few definitions and results concerning the mathematical framework used along the paper, which we have included here for those readers who are not familiar with quantum mechanics. In appendix B we discuss some more results related to the mean-field model.

## 2. A first model

The model we discuss in this section is really an oversimplified toy model of a stock market based on the following assumptions:

1. our market consists of $L$ traders exchanging a single kind of share;
2. the total number of shares, $N$, is fixed in time;
3. a trader can only interact with a single other trader: i.e. the traders feel only a two-body interaction;
4. the traders can only buy or sell one share in any single transaction;
5. there exists an unique price for the share, fixed by the market. In other words, we are not considering any difference between the put and the buy prices;
6. the price of the share changes with discrete steps, multiples of a given monetary unit;
7. each trader has a huge quantities of cash that he can use to buy shares but which does not enter, in the present model, in the definition of his portfolio whose value is fixed only by the number of shares.
Let us briefly comment the above assumptions: of course assuming that there is only a single kind of share may appear rather restrictive but we believe that more species of shares can be introduced without major changes. However, along all this paper we only work in this hypothesis just to simplify the treatment. The third assumption above simply means that it is not possible for, say, the traders $t_{1}, t_{2}$ and $t_{3}$ to interact with each other at the same time: however, $t_{1}$ can interact directly with $t_{3}$ or via its interaction with $t_{2}: t_{1}$ interacts with $t_{2}$ and $t_{2}$ interacts with $t_{3}$. This is a typical simplification in many-body theory where often all the $N$-body interactions, $N \geqslant 3$, are assumed to be negligible with respect to the two-body interaction. Assumptions 4, 5 and 7 are useful to simplify the model and to allow us to extract some driving ideas to construct more realistic models. Finally, as we have seen, assumption 6 is a natural one. Most of these assumptions will be relaxed in the next section.

As we discuss in the appendix, the time behaviour of this model can be described by an operator called the Hamiltonian of the model, which describes the free evolution of the model
plus the effects due to the interaction between the traders. The Hamiltonian of this simple model is the following:
$H=H_{0}+H_{\text {price }}, \quad H_{0}=\sum_{l=1}^{L} \alpha_{l} a_{l}^{\dagger} a_{l}+\sum_{i, j=1}^{L} p_{i j} a_{i} a_{j}^{\dagger}, \quad H_{\text {price }}=\epsilon p^{\dagger} p$,
where the following commutation rules are assumed:

$$
\begin{equation*}
\left[a_{l}, a_{n}^{\dagger}\right]=\delta_{l n} \mathbb{I}, \quad\left[p, p^{\dagger}\right]=\mathbb{I} \tag{2.2}
\end{equation*}
$$

while all the other commutators are zero. The meaning of these operators is discussed in more details in appendix A. Here we just recall that $a_{l}$ and $a_{l}^{\dagger}$, respectively destroys and creates a share in the portfolio of $t_{l}$, while the operators $p$ and $p^{\dagger}$ modify the price of the share: $p$ makes the price decrease of $\epsilon$, while $p^{\dagger}$ makes it increase of the same quantity. The coefficients $p_{i j}$ 's take value 1 or 0 depending on the fact that $t_{i}$ interacts with $t_{j}$ or not. We also assume that $p_{i i}=0$ for all $i$, which simply means that $t_{i}$ does not interact with himself. For those who are familiar with second quantization, there is an easy interpretation for the Hamiltonian above, which can be deduced also from what is discussed in appendix A: while $\epsilon p^{\dagger} p+\sum_{l=1}^{L} \alpha_{l} a_{l}^{\dagger} a_{l}$ describes the free evolution of the operators $\left\{a_{l}\right\}$ and $p$, whose physical meaning will be considered again later on in this section, the single contribution $a_{i} a_{j}^{\dagger}$ of the interaction Hamiltonian $\sum_{i, j=1}^{L} p_{i j} a_{i} a_{j}^{\dagger}$ destroys a share belonging to the trader $t_{i}$ and creates a share in the portfolio of the trader $t_{j}$. In other words: if $p_{i j}=1$ then the trader $t_{i}$ sells a share to $t_{j}$. However, since $H$ must be self-adjoint (for mathematical and physical reasons), then $p_{i j}=1$ also implies $p_{j i}=1$. This means that the interaction Hamiltonian contains both the possibility that $t_{i}$ sells a share to $t_{j}$ and the possibility that $t_{j}$ sells a share to $t_{i}$. Different values of $\alpha_{i}$ in the free Hamiltonian are then used to introduce an ability of the trader, which will make more likely that the most expert trader sells or buys his shares to the other traders so to increase the value of his portfolio.

As we will discuss in appendix A , the time evolution of an operator $X$ of the model is $X(t)=\mathrm{e}^{\mathrm{i} H t} X \mathrm{e}^{-\mathrm{i} H t}$ and it satisfies the following Heisenberg differential equation: $\frac{\mathrm{d} X(t)}{\mathrm{d} t}=\mathrm{i}^{\mathrm{i} H t}[H, X] \mathrm{e}^{-\mathrm{i} H t}=\mathrm{i}[H, X(t)]$. The only observables whose time evolution we are interested in are, clearly, the price of the share and the number of shares of each traders. Indeed, as we have already remarked, within our simplified scheme there is no room for the cash of the trader! The price operator $\hat{P}$ is $\hat{P}=\epsilon p^{\dagger} p$, while the $j$-number operator is $\hat{n}_{j}=a_{j}^{\dagger} a_{j}$, which represents the number of shares that $t_{j}$ possesses. The operator total number of shares is finally $\hat{N}=\sum_{j=1}^{L} \hat{n}_{j}=\sum_{j=1}^{L} a_{j}^{\dagger} a_{j}$. The choice of $H$ in (2.1) is suggested by the requirement (2) above. Indeed it is easy to check, using (2.2), that $[H, \hat{N}]=0$. This implies that the time evolution of $\hat{N}, \hat{N}(t)=\mathrm{e}^{\mathrm{i} H t} \hat{N} \mathrm{e}^{-\mathrm{i} H t}$ is trivial: $\hat{N}(t)=\hat{N}$ for all $t$. However this does not imply also that $\left[H, \hat{n}_{j}\right]=0$, which, as a matter of fact, is not true in general. This is clear from the definition of $H$ : the term $\sum_{l=1}^{L} \alpha_{l} a_{l}^{\dagger} a_{l}$ does not change the number of shares of the different traders, but only counts this number. On the contrary, $\sum_{i, j=1}^{L} p_{i j} a_{i} a_{j}^{\dagger}$ destroys a share belonging to $t_{i}$ but, at the same time, creates another share in the portfolio of the trader $t_{j}$. In this operation, the number of the shares of the single traders are changed, but the total number of shares remains constant! It may be worth noticing that if all the $p_{i j}$ are zero, i.e. if there is no interaction between the traders, then we also get $\left[H, \hat{n}_{j}\right]=0$ : our model produce a completely stationary market, as it is expected.

We implement assumptions (5) and (6) by requiring that the price operator $\hat{P}$ has the form given above, $\hat{P}=\epsilon p^{\dagger} p$, where $\epsilon$ is the monetary unit. Such an operator is assumed to be a part of $H$, see (2.1). Also, because of the simplifications which are assumed in this toy model, $\hat{P}$ is clearly a constant of motion: $[H, \hat{P}]=0$. This is not a realistic assumption, and
will be relaxed in the next sections. However, it is assumed here since it allows us a better understanding of the meaning of the $\alpha_{l}$ 's, as we will discuss later.

In order to describe a state of the system in which at time $t=0$ the portfolio of the first trader consists of $n_{1}$ shares, the one of $t_{2}$ of $n_{2}$ shares, and so on, and the price of the share is $P=M \epsilon$, we should impose that the market is in a vector state $\omega_{n_{1}, n_{2}, \ldots, n_{L} ; M}$, see appendix A, defined by the vector

$$
\begin{equation*}
\varphi_{n_{1}, n_{2}, \ldots, n_{L} ; M}:=\frac{1}{\sqrt{n_{1}!n_{2}!\cdots n_{L}!M!}}\left(a_{1}^{\dagger}\right)^{n_{1}}\left(a_{2}^{\dagger}\right)^{n_{2}} \cdots\left(a_{L}^{\dagger}\right)^{n_{L}}\left(p^{\dagger}\right)^{M} \varphi_{0}, \tag{2.3}
\end{equation*}
$$

where $\varphi_{0}$ is the vacuum of the model: $a_{j} \varphi_{0}=p \varphi_{0}=0$ for all $j=1,2, \ldots, L$. If $X \in \mathfrak{A}, \mathfrak{A}$ being the algebra of the observables of our market, then we put

$$
\begin{equation*}
\omega_{n_{1}, n_{2}, \ldots, n_{L} ; M}(X)=\left\langle\varphi_{n_{1}, n_{2}, \ldots, n_{L} ; M}, X \varphi_{n_{1}, n_{2}, \ldots, n_{L} ; M}\right\rangle, \tag{2.4}
\end{equation*}
$$

and $\langle$,$\rangle is the scalar product in the Hilbert space of the theory, see again the appendix.$ The Heisenberg equations of motion (A.2) for the annihilation operators $a_{l}(t)$ produce the following very simple differential equation:

$$
\begin{equation*}
\mathrm{i} \dot{a}(t)=X a(t) \tag{2.5}
\end{equation*}
$$

where we have introduced the matrix $X$, independent of time, and the vector $a(t)$ as follows:

$$
X \equiv\left(\begin{array}{ccccccc}
\alpha_{1} & p_{21} & p_{31} & \cdot & \cdot & p_{L-11} & p_{L 1} \\
p_{12} & \alpha_{2} & p_{32} & \cdot & \cdot & \cdot & p_{L 2} \\
p_{13} & p_{23} & \alpha_{3} & \cdot & \cdot & \cdot & \cdot \\
\cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\
\cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\
p_{1 L-1} & p_{2 L-1} & p_{3 L-1} & \cdot & \cdot & \alpha_{L-1} & p_{L L-1} \\
p_{1 L} & p_{2 L} & p_{3 L} & \cdot & \cdot & p_{L-1 L} & \alpha_{L}
\end{array}\right), \quad a(t) \equiv\left(\begin{array}{c}
a_{1}(t) \\
a_{2}(t) \\
a_{3}(t) \\
\cdot \\
\cdot \\
a_{L-1}(t) \\
a_{L}(t)
\end{array}\right)
$$

Note that, due to the conditions on the $p_{i j}$ 's, and since all the $\alpha_{l}$ 's are real, the matrix $X$ is self-adjoint. Equation (2.5) can now be solved as follows: let $V$ be the (unitary) matrix which diagonalizes $X: V^{\dagger} X V=\operatorname{diag}\left\{x_{1}, x_{2}, \ldots, x_{L}\right\}=: X_{d}, x_{j}$, being its eigenvalues, $j=1,2, \ldots, L$. Note that, of course, $V$ does not depend on time. Then, putting

$$
U(t)=\left(\begin{array}{ccccccc}
\mathrm{e}^{\mathrm{i} x_{1} t} & 0 & 0 & . & . & . & 0 \\
0 & \mathrm{e}^{\mathrm{i} x_{2} t} & 0 & . & . & . & 0 \\
0 & 0 & \mathrm{e}^{\mathrm{i} x_{3} t} & . & . & . & 0 \\
. & . & . & . & . & . & . \\
. & . & . & . & . & . & \cdot \\
0 & 0 & 0 & . & . & . & \mathrm{e}^{\mathrm{i} x_{L} t}
\end{array}\right)
$$

we get

$$
\begin{equation*}
a(t)=V U(t) V^{\dagger} a(0), \tag{2.6}
\end{equation*}
$$

where, as it is clear, $a(0)^{T}=\left(a_{1}, a_{2}, \ldots, a_{L}\right)$. If we further introduce the adjoint of the vector $a(t), a^{\dagger}(t)=\left(a_{1}^{\dagger}(t), a_{2}^{\dagger}(t), \ldots, a_{L}^{\dagger}(t)\right)=a^{\dagger}(0) V U^{\dagger}(t) V^{\dagger}$, we can explicitly check that $\hat{N}$ is a constant of motion. Indeed we have

$$
\begin{aligned}
\hat{N}(t) & =a_{1}^{\dagger}(t) a_{1}(t)+a_{2}^{\dagger}(t) a_{2}(t)+\cdots+a_{L}^{\dagger}(t) a_{L}(t)=a^{\dagger}(t) \cdot a(t) \\
& =\left(a^{\dagger}(0) V U^{\dagger}(t) V^{\dagger}\right) \cdot\left(V U(t) V^{\dagger} a(0)\right)=a^{\dagger}(0) \cdot a^{\dagger}(0)=\hat{N}(0) .
\end{aligned}
$$

In order to analyse the time behaviour of the different $\hat{n}_{j}(t)$, we simply have to compute the mean value $n_{j}(t)=\omega_{n_{1}, n_{2}, \ldots, n_{L} ; M}\left(\hat{n}_{j}(t)\right)$. This means that, for $t=0$, the first trader


Figure 1. $n_{3}(t)$ for $\alpha_{3}=3$ (left), $\alpha_{3}=10$ (middle), $\alpha_{3}=100$ (right).
possesses $n_{1}$ shares, the second trader possesses $n_{2}$ shares, and so on, and that the price of the share is $M \epsilon$. It should be mentioned that the only way in which a matrix element like $\omega_{n_{1}, n_{2}, \ldots, n_{L} ; M}\left(a_{j}^{k}\left(a_{l}^{\dagger}\right)^{m}\right)$ can be different from zero is when $j=l$ and $k=m$. This follows from the orthonormality of the set $\left\{\varphi_{n_{1}, n_{2}, \ldots, n_{L} ; M}\right\}$ which is a direct consequence of the canonical commutation relations.

The easiest way to get the analytic expression for $n_{j}(t)$ is to fix the number of the traders, starting with the simplest situation: $L=2$. In this case we find that

$$
\left\{\begin{array}{l}
n_{1}(t)=\frac{1}{\Omega^{2}}\left\{n_{1}\left(\alpha^{2}+2 p^{2}(1+\cos (\Omega t))\right)+2 p^{2} n_{2}(1-\cos (\Omega t))\right\}  \tag{2.7}\\
\left.n_{2}(t)=\frac{2 p^{2} n_{1}}{\Omega^{2}}(1-\cos (\Omega t))\right)+n_{2}\left(1+\frac{2 p^{2}}{\Omega^{2}}(\cos (\Omega t)-1)\right)
\end{array}\right.
$$

where we have defined $\Omega^{2}=\alpha^{2}+4 p^{2}$, with $\alpha=\alpha_{2}-\alpha_{1}$ and $p=p_{12}=p_{21}$.
It is not hard to check that $n_{1}(t)+n_{2}(t)=n_{1}+n_{2}$, as expected. Also, if $p=0$ then we find $n_{1}(t)=n_{1}$ and $n_{2}(t)=n_{2}$ for all $t$. This is natural and expected, since when $p=0$ there is no interaction at all between the traders, so that there is no reason for $n_{1}(t)$ and $n_{2}(t)$ to change in time. Another interesting consequence of (2.7) is that, if $n_{1}=n_{2}=n$, that is if the two traders start with the same number of shares, they do not change this equilibrium during the time: we find again $n_{1}(t)=n_{2}(t)=n$. Also this result is expected, since both $t_{1}$ and $t_{2}$ possess the same amount of money (their huge sources!) and the same number of shares. The role of $\alpha_{1}$ and $\alpha_{2}$, in this case, is unessential. It is further clear that $n_{j}(t)$ is a periodic function whose period, $T=\frac{2 \pi}{\Omega}$, decreases for $|\alpha|=\left|\alpha_{1}-\alpha_{2}\right|$ and $p$ increasing. Finally, if we call $\Delta n_{j}=\max _{t \in[0, T]}\left|n_{j}(t)-n_{j}(0)\right|$, which represents the highest variation of $n_{j}(t)$ in a period, we can easily check that $\Delta n_{j}$ increases when $\left|n_{1}(0)-n_{2}(0)\right|$ increases and when $\Omega$ decreases.

Remark. It is worth remarking that, since the number of shares should be integer, while the functions $n_{1}(t)$ and $n_{2}(t)$ are not integers for general values of $t$, we could introduce a sort of time per the mth transaction, $\tau_{m}$, chosen in such a way that $n_{j}\left(\tau_{1}\right), n_{j}\left(\tau_{2}\right), \ldots$ are all integers, $j=1,2$.

Let us now consider a market with three traders. In figure 1 we plot $n_{3}(t)$ with the initial conditions $n_{1}=40, n_{2}=n_{3}=0$, with $p_{12}=p_{13}=p_{23}=1$ and different values of $\alpha_{1}, \alpha_{2}$ and $\alpha_{3}$. In the figure on the left we have $\left(\alpha_{1}, \alpha_{2}, \alpha_{3}\right)=(1,2,3)$, in the one in the middle $\left(\alpha_{1}, \alpha_{2}, \alpha_{3}\right)=(1,2,10)$ and in the one in the right $\left(\alpha_{1}, \alpha_{2}, \alpha_{3}\right)=(1,2,100)$

This plot, together with many others which can be obtained, e.g., considering different initial conditions, suggests to interpret $\alpha_{j}$ as a sort of inertia: the larger the value of $\alpha_{j}$, the bigger the tendency of $t_{j}$ of keeping the number of his shares constant in time! We could also think of $\alpha_{j}^{-1}$ as a sort of information reaching $t_{j}$ (but not the other traders): if $\alpha_{j}$ is large then


Figure 2. $n_{1}(t), n_{2}(t), n_{3}(t)$ (first row) and $n_{4}(t), n_{5}(t)$ (second row) for $\alpha_{j}$ and $n_{j}$ as above.
not much information reaches $t_{j}$ which has therefore no input to optimize his interaction with the other traders.

In this case, and also for more traders, it is not evident from our plots if a periodic behaviour is again recovered. In any case, at least a quasi-periodic behaviour is observed with a quasi-period which is compatible with the same $T$ found in the case of the two traders. It is clear, however, that the analytic expression of $n_{j}(t)$ can be found using (2.6).

As for the $L=2$ situation, we recover that if $n_{1}=n_{2}=n_{3}=n$, then $n_{1}(t)=n_{2}(t)=$ $n_{3}(t)=n$, for all $t$. Moreover, if $n_{1} \simeq n_{2} \simeq n_{3} \simeq n$, then $n_{j}(t)$ have all small oscillations around $n$. But, if $n_{1} \simeq n_{2} \neq n_{3}$, and if $p_{i j}=1$ for all $i, j$ with $i \neq j$, then all the functions $n_{j}(t)$ change considerably with time. The reason is the following: since $n_{3}$ differs from $n_{1}$ and $n_{2}$, it is natural to expect that $n_{3}(t)$ changes with time. But, since $N=n_{1}(t)+n_{2}(t)+n_{3}(t)$ must be constant, both $n_{2}(t)$ and $n_{1}(t)$ must change in time as well. The same conclusion can be deduced also if $p_{23}=0$ while all the other $p_{i j}$ 's are equal to 1 : even if $t_{2}$ does not interact with $t_{3}$, the fact that $t_{1}$ interacts with both $t_{2}$ and $t_{3}$, together with the fact that $N$ must be constant, implies again that all the $n_{j}(t)$ 's need to change in time. Finally, it is clear that if $p_{13}=p_{23}=0$, then $t_{3}$ interact neither with $t_{1}$ nor with $t_{2}$ and, indeed, we find that $n_{3}(t)$ does not change with time: this is a consequence of the fact that, in this case, $\left[H, \hat{n}_{3}\right]=0$.

Analogous conclusions can be deduced also for five (or more) traders. In particular figure 2 shows that there is no need for all the traders to interact among them to have a non-trivial time behaviour. Indeed, even if $p_{15}=p_{25}=0$, which means that $t_{5}$ may only interact directly with $t_{3}$ and $t_{4}$, we get the following plots for $\left(\alpha_{1}, \alpha_{2}, \alpha_{3}, \alpha_{4}, \alpha_{5}\right)=(1,2,3,4,5)$ and $\left(n_{1}, n_{2}, n_{3}, n_{4}, n_{5}\right)=(40,0,0,0,0)$. We see that the number of shares of each trader changes in time with the same order of magnitude.

We end this section stressing that the interpretation of $\alpha_{j}$ as a sort of inertia is also suggested by the analysis of this larger number of traders.

## 3. A different model

We consider now another model, which differs from the previous one since the cash, the price of the share and the supply of the market are introduced in a non-trivial way. In particular
we require that assumptions $1,2,3,4$ and 6 of the previous section still hold. Moreover we require the following:
(a) When the tendency of the market to buy a share, i.e. the market demand, increases then the price of the share increases as well. Equivalently, when the tendency of the market to sell a share, i.e. the market supply, increases then the price of the share decreases.
(b) For our convenience the demand and the supply are expressed in term of natural numbers.
(c) We take $\epsilon=1$ in the following: 1 is therefore the unit of money.

The formal Hamiltonian of the model is the following operator:

$$
\left\{\begin{array}{l}
\tilde{H}=H_{0}+\tilde{H}_{I}, \quad \text { where }  \tag{3.1}\\
H_{0}=\sum_{l=1}^{L} \alpha_{l} a_{l}^{\dagger} a_{l}+\sum_{l=1}^{L} \beta_{l} c_{l}^{\dagger} c_{l}+o^{\dagger} o+p^{\dagger} p \\
\tilde{H}_{I}=\sum_{i, j=1}^{L} p_{i j}\left(a_{i}^{\dagger} a_{j}\left(c_{i} c_{j}^{\dagger}\right)^{\hat{P}}+a_{i} a_{j}^{\dagger}\left(c_{j} c_{i}^{\dagger}\right)^{\hat{P}}\right)+\left(o^{\dagger} p+p^{\dagger} o\right)
\end{array}\right.
$$

where, as before, $\hat{P}=p^{\dagger} p$. Here the following commutation rules are assumed:

$$
\begin{equation*}
\left[a_{l}, a_{n}^{\dagger}\right]=\left[c_{l}, c_{n}^{\dagger}\right]=\delta_{l n} \mathbb{I}, \quad\left[p, p^{\dagger}\right]=\left[o, o^{\dagger}\right]=\mathbb{I} \tag{3.2}
\end{equation*}
$$

while all the other commutators are zero. As for the previous model we assume that $p_{i i}=0$. Here the operators $a_{l}^{\sharp}$ and $p^{\sharp}$ have the same meaning as in the previous section, while $c_{l}^{\sharp}$ and $o^{\sharp}$ are respectively the cash and the supply operators. The states in (2.4) must be replaced by the states

$$
\begin{equation*}
\omega_{\{n\} ;\{k\} ; O ; M}(.)=\left\langle\varphi_{\{n\} ;\{k\} ; O ; M}, . \varphi_{\{n\} ;\{k\} ; O ; M}\right\rangle, \tag{3.3}
\end{equation*}
$$

where $\{n\}=n_{1}, n_{2}, \ldots, n_{L},\{k\}=k_{1}, k_{2}, \ldots, k_{L}$ and

$$
\begin{equation*}
\varphi_{\{n\} ;\{k\} ; O ; M}:=\frac{\left(a_{1}^{\dagger}\right)^{n_{1}} \cdots\left(a_{L}^{\dagger}\right)^{n_{L}}\left(c_{1}^{\dagger}\right)^{k_{1}} \cdots\left(c_{L}^{\dagger}\right)^{k_{L}}\left(o^{\dagger}\right)^{o}\left(p^{\dagger}\right)^{M}}{\sqrt{n_{1}!\ldots n_{L}!k_{1}!\ldots k_{L}!O!M!}} \varphi_{0} \tag{3.4}
\end{equation*}
$$

Here $\varphi_{0}$ is the vacuum of the model: $a_{j} \varphi_{0}=c_{j} \varphi_{0}=p \varphi_{0}=o \varphi_{0}=0$, for $j=1,2, \ldots, L$.
Let us now see what is the meaning of the Hamiltonian above and for which reason we call it formal.
$H_{0}$ contains all that is related to the free dynamics of the model.
$\tilde{H}_{I}$ is the interaction Hamiltonian, whose terms have a natural interpretation as follows:
The presence of $o^{\dagger} p$ implies that when the supply increases then the price must decrease. Of course $p^{\dagger} o$ produces exactly the opposite effect.

The presence of $a_{i}^{\dagger} a_{j}\left(c_{i} c_{j}^{\dagger}\right)^{\hat{P}}$ implies that $t_{i}$ increases of one unit the number of shares in his portfolio but, at the same time, his cash decreases because of $c_{i}^{\hat{P}}$, that is it must decrease of as many units of cash as the price operator $\hat{P}$ demands. Moreover, the trader $t_{j}$ behaves exactly in the opposite way: he has one share less because of $a_{j}$ but his cash increases because of $\left(c_{j}^{\dagger}\right)^{\hat{P}}$. Of course, the Hermitian conjugate term $a_{i} a_{j}^{\dagger}\left(c_{j} c_{i}^{\dagger}\right)^{\hat{P}}$ in $\tilde{H}_{I}$ produces a specular effect for the two traders.

As in the previous section, if $\tilde{H}_{I}=0$, then there is no nontrivial dynamics of the relevant observables of the system, like $c_{j}^{\dagger} c_{j}$ and $n_{j}^{\dagger} n_{j}$. This can also be seen as a criterium to fix the free Hamiltonian of the system: it is only the interaction between the traders which may modify their status!

However, despite of this clear physical interpretation, the Hamiltonian in (3.1) suffers of a technical problem: since $c_{j}$ and $c_{j}^{\dagger}$ are not self-adjoint operators, it is not obvious how to define, for instance, the operator $c_{j}^{\hat{P}}$. Indeed, if we formally write $c_{j}^{\hat{P}}$ as $\mathrm{e}^{\hat{P} \log c_{j}}$, then we cannot use functional calculus to define $\log c_{j}$. Also, we cannot use a simple series expansion
since the operators involved are all unbounded so that the series we get is surely not norm convergent and many domain problems appear. For this reason, we find more convenient to replace $\tilde{H}$ with an effective Hamiltonian, $H$, defined as

$$
\left\{\begin{array}{l}
H=H_{0}+H_{I}, \quad \text { where }  \tag{3.5}\\
H_{0}=\sum_{l=1}^{L} \alpha_{l} a_{l}^{\dagger} a_{l}+\sum_{l=1}^{L} \beta_{l} c_{l}^{\dagger} c_{l}+o^{\dagger} o+p^{\dagger} p \\
H_{I}=\sum_{i, j=1}^{L} p_{i j}\left(a_{i}^{\dagger} a_{j}\left(c_{i} c_{j}^{\dagger}\right)^{M}+a_{i} a_{j}^{\dagger}\left(c_{j} c_{i}^{\dagger}\right)^{M}\right)+\left(o^{\dagger} p+p^{\dagger} o\right)
\end{array}\right.
$$

where $M=\omega_{\{n\} ;\{k\} ; O ; M}(\hat{P})$. Note that, because of the fact that $p_{i i}=0$, there is no difference in $H_{I}$ above if we write $\left(c_{i} c_{j}^{\dagger}\right)^{M}$ or $\left(c_{i}\right)^{M}\left(c_{j}^{\dagger}\right)^{M}$ even if the two operator $c_{j}$ and $c_{j}^{\dagger}$ do not commute. Note also that if we consider a state $\omega$ over $\mathfrak{A}$ different from $\omega_{\{n\} ;\{k\} ; O ; M}$, as we will do in the next section, then $\omega(\hat{P})$ could be different from $M$.

Three integrals of motion for our model trivially exist:

$$
\begin{equation*}
\hat{N}=\sum_{i=1}^{L} a_{i}^{\dagger} a_{i}, \quad \hat{K}=\sum_{i=1}^{L} c_{i}^{\dagger} c_{i} \quad \text { and } \quad \hat{\Gamma}=o^{\dagger} o+p^{\dagger} p \tag{3.6}
\end{equation*}
$$

This can be easily checked since the canonical commutation relations in (3.2) imply that $[H, \hat{N}]=[H, \hat{\Gamma}]=[H, \hat{K}]=0$.

The fact that $\hat{N}$ is conserved clearly means that no new share is introduced in the market. Of course, also the total amount of money must be a constant of motion since the cash is assumed to be used only to buy shares. Since also $\hat{\Gamma}$ commutes with $H$, moreover, if the mean value of $o^{\dagger} o$ increases with time then necessarily the mean value of the price operator must decrease and vice versa. This is exactly the mechanism assumed in point (a) at the beginning of this section.

Remark. It may be worth noticing that this is not the only way in which requirement (a) could be implemented, but it is surely the simplest one. Just to give few other examples, we could ask for one the following combinations to remain constant in time: $\left(o^{\dagger} o\right)^{2}+\left(p^{\dagger} p\right)^{2}, o^{\dagger} o p^{\dagger} p$ or many others.

Another consequence of the definition of $H$ is that $L$ other constants of motion also exist. They are the following operators:

$$
\begin{equation*}
\hat{Q}_{j}=a_{j}^{\dagger} a_{j}+\frac{1}{M} c_{j}^{\dagger} c_{j}, \tag{3.7}
\end{equation*}
$$

for $j=1,2, \ldots, L$. This can be checked explicitly computing [ $H, \hat{Q}_{j}$ ] and proving that all these commutators are zero. But we can also understand this feature simply noticing that (i) $\hat{Q}_{j}$ commutes trivially with $H_{0}$ and (ii) the term $a_{i}^{\dagger} a_{j}\left(c_{i} c_{j}^{\dagger}\right)^{M}$ in $H_{I}$ obviously preserves not only the total number of shares and the total amount of cash, but also a certain combination of the shares and the cash: as far as $t_{i}$ is concerned, $a_{i}^{\dagger}$ increases of one unit the number of shares while $c_{i}^{M}$ decreases of $M$ units the amount of cash. This means that if a certain vector $\Psi$ represents $n_{i}$ shares and $k_{i}$ units of cash, then $a_{i}^{\dagger} c_{i}^{M} \Psi$ describes $n_{i}+1$ shares and $k_{i}-M$ units of cash. Therefore we have $\hat{Q}_{i} \Psi=\left(n_{i}+\frac{1}{M} k_{i}\right) \Psi$ and $\hat{Q}_{i}\left(a_{i}^{\dagger} c_{i}^{M} \Psi\right)=\left(n_{i}+1+\frac{1}{M}\left(k_{i}-M\right)\right)\left(a_{i}^{\dagger} c_{i}^{M} \Psi\right)=\left(n_{i}+\frac{1}{M} k_{i}\right)\left(a_{i}^{\dagger} c_{i}^{M} \Psi\right)$. So, it is not surprising that $\left[\hat{Q}_{i}, a_{i}^{\dagger} c_{i}^{M}\right]=0$ and, as a consequence, that $\left[\hat{Q}_{i}, H\right]=0$.

The Hamiltonian (3.5) contains a contribution, $h_{p o}=o^{\dagger} o+p^{\dagger} p+\left(o^{\dagger} p+p^{\dagger} o\right)$, which is decoupled from the other terms. This means that, within our model, the time evolution of the
supply and the price operators do not depend on the number of shares or on the cash, and can be deduced referring only to $h_{p o}$. The Heisenberg equations of motion are the following:

$$
\left\{\begin{array}{l}
\mathrm{i} \dot{o}(t)=o(t)+p(t)  \tag{3.8}\\
\mathrm{i} \dot{p}(t)=o(t)+p(t),
\end{array}\right.
$$

which shows that $o(t)-p(t)$ is constant in $t$, so that $\hat{\Delta}=o-p$ is still another integral of motion. Solving this system we get $o(t)=\frac{1}{2}\left\{o\left(\mathrm{e}^{-2 \mathrm{it}}+1\right)+p\left(\mathrm{e}^{-2 \mathrm{i} t}-1\right)\right\}$ and $p(t)=\frac{1}{2}\left\{p\left(\mathrm{e}^{-2 \mathrm{i} t}+1\right)+o\left(\mathrm{e}^{-2 \mathrm{i} t}-1\right)\right\}$. It is now trivial to check explicitly that both $\hat{\Delta}(t)=o(t)-p(t)$ and $\hat{\Gamma}(t)=o^{\dagger}(t) o(t)+p^{\dagger}(t) p(t)$ do not depend on time. If we now compute the mean value of the price and supply operators on a state number we get

$$
\left\{\begin{array}{l}
P_{r}(t)=\frac{1}{2}\left\{P_{r}+O_{f}+\left(P_{r}-O_{f}\right) \cos (2 t)\right\}  \tag{3.9}\\
O_{f}(t)=\frac{1}{2}\left\{P_{r}+O_{f}-\left(P_{r}-O_{f}\right) \cos (2 t)\right\},
\end{array}\right.
$$

where we have called $O_{f}(t)=\omega_{\{n\} ;\{k\} ; O ; M}\left(o^{\dagger}(t) o(t)\right)$ and $P_{r}(t)=\omega_{\{n\} ;\{k\} ; O ; M}\left(p^{\dagger}(t) p(t)\right)$. Recall that $P_{r}=P_{r}(0)=M$. Equations (3.9) show that, if $O_{f}=P_{r}$ then $O_{f}(t)=P_{r}(t)=$ $O_{f}$ for all $t$ while, if $O_{f} \simeq P_{r}$ then $O_{f}(t)$ and $P_{r}(t)$ are almost constant. In the following we will replace $P_{r}(t)$ with an integer value, the value $M$ which appears in the Hamiltonian (3.5), which is therefore fixed after the solution (3.9) is found. This value is obtained by taking a suitable mean of $P_{r}(t)$ or working in one of the following assumptions: (i) $O_{f}=P_{r}$; or (ii) $O_{f} \simeq P_{r}$ or yet (iii) $\left|O_{f}+P_{r}\right| \gg\left|P_{r}-O_{f}\right|$. In these last two situations we may replace $P_{r}(t)$, with a temporal mean, $\left\langle P_{r}(t)\right\rangle$, since there is not much difference between these two quantities.

Let us now recall that the main aim of each trader is to improve the total value of his portfolio, which we define as follows:

$$
\begin{equation*}
\hat{\Pi}_{j}(t)=\gamma \hat{n}_{j}(t)+\hat{k}_{j}(t) \tag{3.10}
\end{equation*}
$$

Here we have introduced the value of the share $\gamma$ as decided by the market, which does not necessarily coincides with the amount of money which is paid to buy the share. As is clear, $\hat{\Pi}_{j}(t)$ is the sum of the complete value of the shares, plus the cash. The fact that for each $j$ the operator $Q_{j}$ is an integral of motion allows us to rewrite the operator $\hat{\Pi}_{j}(t)$ only in terms of $\hat{n}_{j}(t)$ and of the initial conditions. We find

$$
\begin{equation*}
\hat{\Pi}_{j}(t)=\hat{\Pi}_{j}(0)+(\gamma-M)\left(\hat{n}_{j}(t)-\hat{n}_{j}(0)\right) \tag{3.11}
\end{equation*}
$$

In order to get the time behaviour of the portfolio, it is enough to obtain $\hat{n}_{j}(t)$. If we write the Heisenberg equation for $\hat{n}_{j}(t), \dot{\hat{n}}_{j}(t)=\mathrm{i}\left[H, \hat{n}_{j}(t)\right]$, we see that this equation involves the time evolution of $a_{j}, c_{j}$ and their adjoint. The equations of motion for these operators should be added to close the system, and the final system of differential equations cannot be solved exactly. The easiest way to proceed is to develop the following simple perturbative expansion, well known in quantum mechanics:
$\hat{n}_{j}(t)=\mathrm{e}^{\mathrm{i} H t} \hat{n}_{j} \mathrm{e}^{-\mathrm{i} H t}=\hat{n}_{j}+\mathrm{i} t\left[H, \hat{n}_{j}\right]+\frac{(\mathrm{i} t)^{2}}{2!}\left[H, \hat{n}_{j}\right]_{2}+\frac{(\mathrm{i} t)^{3}}{3!}\left[H, \hat{n}_{j}\right]_{3}+\cdots$,
where $\left[H, \hat{n}_{j}\right]_{1}=\left[H, \hat{n}_{j}\right]=H \hat{n}_{j}-\hat{n}_{j} H$ and $\left[H, \hat{n}_{j}\right]_{n+1}=\left[H,\left[H, \hat{n}_{j}\right]_{n}\right]$ for $n \geqslant 1$, and then to take its mean value on a state $\omega_{\{n\} ;\{k\} ; O ; M}$ up to the desired order of accuracy. Of course, we can compute as many contributions of the above expansion as we want. However, the expression for $\left[H, \hat{n}_{j}\right]_{n}$ becomes more and more involved as $n$ and $L$ increase. Just as an example, we consider here the case $L=2$ : up to the third order in time we find

$$
\begin{equation*}
n_{1}(t)=\omega_{\{n\} ;\{k\} ; O ; M}\left(\hat{n}_{1}(t)\right)=n_{1}+t^{2} p_{12}^{2}\left(\epsilon_{+}^{2}-\epsilon_{-}^{2}\right)+O\left(t^{4}\right) \tag{3.13}
\end{equation*}
$$

where $\epsilon_{ \pm}$are related to the state $\omega_{\{n\} ;\{k ; ; O ; M}$ as follows:
$\epsilon_{+}=\sqrt{\left(n_{1}+1\right) n_{2} \frac{k_{1}!}{\left(k_{1}-M\right)!} \frac{\left(k_{2}+M\right)!}{k_{2}!}}, \quad \epsilon_{-}=\sqrt{\left(n_{2}+1\right) n_{1} \frac{\left(k_{1}+M\right)!}{k_{1}!} \frac{k_{2}!}{\left(k_{2}-M\right)!}}$.
Of course, in order to have all the above quantities well defined, we need to have both $k_{2}-M \geqslant 0$ and $k_{1}-M \geqslant 0$. This is a natural requirement since it simply states that a trader can buy a share only if he has the money to pay for it!

Remarks. (1) This solution has some analogies with that given in (2.7). Indeed, if we expand $n_{1}(t)$ in (2.7) as a power of $t$, we find an expression which is very close to equation (3.13). In particular, we find that for both models there is no contribution coming from $\alpha_{j}$ (and from $\beta_{j}$ here) up to the order $t^{3}$. Also, for this model, if $t_{1}$ and $t_{2}$ possess the same amount of cash for $t=0, k_{1}=k_{2}$, then, since $\epsilon_{+}^{2}-\epsilon_{-}^{2}$ turns out to be proportional to $n_{2}-n_{1}$, we deduce that $n_{1}(t)=n_{2}(t)$ if $n_{1}=n_{2}$. This is again exactly the same conclusion we have obtained in section $2: n_{1}=n_{2}$ is a stability condition.
(2) Using the fact that $Q_{1}$ is constant we can also find the value of the cash of $t_{1}$ as a function of time: $k_{1}(t)=k_{1}-M t^{2} p_{12}^{2}\left(\epsilon_{+}^{2}-\epsilon_{-}^{2}\right)+O\left(t^{4}\right)$ while its portfolio evolves like

$$
\begin{equation*}
\Pi_{1}(t)=\Pi_{1}(0)+(\gamma-M) t^{2} p_{12}^{2}\left(\epsilon_{+}^{2}-\epsilon_{-}^{2}\right)+O\left(t^{4}\right) \tag{3.14}
\end{equation*}
$$

(3) This formula allows us to get some conclusions concerning the time behaviour of the portfolio of $t_{1}$ for small time. In particular we can deduce that
if $k_{1}=k_{2}$ and $n_{1}=n_{2}$ then $k_{j}(t)=k_{j}$ and $n_{j}(t)=n_{j}, j=1,2$. The two traders are already in an equilibrium state and there is no way to let them change their state;
if $k_{1}=k_{2}=: k$ but $n_{1} \neq n_{2}$ then, since $\epsilon_{+}^{2}-\epsilon_{-}^{2}=\frac{(k+M)!}{(k-M)!}\left(n_{2}-n_{1}\right)$, it follows that $\epsilon_{+}^{2}-\epsilon_{-}^{2}>0$ if $n_{2}>n_{1}$ and it is negative otherwise. This implies that, for small $t, n_{1}(t)$ increases with $t$ if $n_{2}>n_{1}$ and decreases if $n_{2}<n_{1}$. This means that the trader with more shares tends to sell some of his shares to the other trader, to increase his liquidity. Moreover, since $k_{1}(t)=Q_{1}-n_{1}(t), k_{1}(t)$ decreases when $n_{1}(t)$ increases and vice versa. We also find

$$
\begin{equation*}
\Pi_{1}(t) \simeq \Pi_{1}(0)+(\gamma-M) t^{2} p_{12}^{2} \frac{(k+M)!}{(k-M)!}\left(n_{2}-n_{1}\right) \tag{3.15}
\end{equation*}
$$

which shows that, if $\gamma>M, \Pi_{1}(t)$ increases with $t$ if $n_{2}>n_{1}$ and decreases if $n_{1}>n_{2}$. This can be understood as follows: if $\gamma>M$, then the market is giving to the shares a larger value than the amount of cash used to buy them. Therefore, if $n_{2}>n_{1}$, since as we have seen $n_{1}(t)$ increases for small $t$, the first trader is paying $M$ for a share whose value is $\gamma>M$. That is why the value of his portfolio increases!

Let now take $n_{1}=n_{2}=: n$ and $k_{1} \neq k_{2}$. In this case after few algebraic computations we see that $\epsilon_{+}^{2}-\epsilon_{-}^{2}>0$ if $k_{1}>k_{2}$ while $\epsilon_{+}^{2}-\epsilon_{-}^{2}<0$ if $k_{1}<k_{2}$. This implies that, if $k_{1}>k_{2}$, then $n_{1}(t)$ increases while $k_{1}(t)$ decreases as $t$ increases. Moreover, if $\gamma>M$, then $\Pi_{1}(t)$ increases its value with $t$. This can be understood again as before: since $n_{1}(t)$ is an increasing function, for $t>0$ small enough, and since the market price of the share $\gamma$ is larger than $M$, $t_{1}$ improves his portfolio since he spends $M$ to get $\gamma$.
(4) As for the second trader, we can easily find $n_{2}(t), k_{2}(t)$ and $\Pi_{2}(t)$ simply recalling that $N=n_{1}(t)+n_{2}(t)$ is constant in time.
(5) The case in which $\gamma<M$ can be analysed in the very same way as before.

## 4. Mean-field approximation

It is clear that the results of the previous section suffer of the two major approximations: first of all our final considerations have been obtained only in the case of two traders. Considering
more traders is technically much harder and goes beyond the real aims of this paper. Second, the perturbation expansion we have introduced in (3.12) gives only an approximated version of the exact solution. In this section we propose a particular version of the model considered before which, under a sufficiently general assumption on $\alpha_{j}$ and $\beta_{j}$, can be explicitly solved in the so-called mean-field approximation. This different version of our model is relevant since it is related to a market in which the number of traders is very large, virtually divergent. In other words, while in the previous section we have considered a stock market with very few traders, using the mean-field approximation we will be able to analyse a different market, namely one with a very large number of traders.

Our model is defined by the same Hamiltonian as in (3.5) but with $M=1$. This is not a major requirement since it corresponds to a renormalization of the price of the share, which we take equal to 1 : if you buy a share, then your liquidity decreases by one unit while it increases, again by one unit, if you sell a share. It is clear that all the same integrals of motion as before exist: $\hat{N}, \hat{K}, \hat{\Gamma}, \hat{\Delta}$ and $Q_{j}=\hat{n}_{j}+\hat{k}_{j}, j=1,2, \ldots, L$. They all commute with $H$, which we now write as

$$
\left\{\begin{array}{l}
H=h+h_{p o}, \quad \text { where }  \tag{4.1}\\
h=\sum_{l=1}^{L} \alpha_{l} \hat{n}_{l}+\sum_{l=1}^{L} \beta_{l} \hat{k}_{l}+\sum_{i, j=1}^{L} p_{i j}\left(a_{i}^{\dagger} a_{j} c_{i} c_{j}^{\dagger}+a_{i} a_{j}^{\dagger} c_{j} c_{i}^{\dagger}\right) \\
h_{p o}=o^{\dagger} o+p^{\dagger} p+\left(o^{\dagger} p+p^{\dagger} o\right)
\end{array}\right.
$$

For $h_{p o}$ we can repeat the same argument as in the previous section and an explicit solution can be found which is completely independent of $h$. In particular we have $\omega_{\{n\} ;\{k\} ; O ; M}(\hat{P})=1$. For this reason, from now on, we will identify $H$ only with $h$ and work only with this Hamiltonian. Let us introduce the operators

$$
\begin{equation*}
X_{i}=a_{i} c_{i}^{\dagger} \tag{4.2}
\end{equation*}
$$

$i=1,2, \ldots, L$. The Hamiltonian $h$ can be rewritten as

$$
\begin{equation*}
h=\sum_{l=1}^{L}\left(\alpha_{l} \hat{n}_{l}+\beta_{l} \hat{k}_{l}\right)+\sum_{i, j=1}^{L} p_{i j}\left(X_{i}^{\dagger} X_{j}+X_{j}^{\dagger} X_{i}\right) \tag{4.3}
\end{equation*}
$$

The following commutation relations hold:
$\left[X_{i}, X_{j}^{\dagger}\right]=\delta_{i j}\left(\hat{k}_{i}-\hat{n}_{i}\right), \quad\left[X_{i}, \hat{n}_{j}\right]=\delta_{i j} X_{i} \quad\left[X_{i}, \hat{k}_{j}\right]=-\delta_{i j} X_{i}$,
which show how the operators $\left\{\left\{X_{i}, X_{i}^{\dagger}, \hat{n}_{i}, \hat{k}_{i}\right\}, i=1,2, \ldots, L\right\}$ are closed under commutation relations. This is quite important, since it produces the following system of differential equations:

$$
\left\{\begin{array}{l}
\dot{X}_{l}=\mathrm{i}\left(\beta_{l}-\alpha_{l}\right) X_{l}+2 \mathrm{i} X_{l}^{(L)}\left(\hat{n}_{l}-\hat{k}_{l}\right) \\
\dot{\hat{n}}_{l}=2 \mathrm{i}\left(X_{l} X_{l}^{(L)^{\dagger}}-X_{l}^{(L)} X_{l}^{\dagger}\right) \\
\hat{k}_{l}=-2 \mathrm{i}\left(X_{l} X_{l}^{(L)^{\dagger}}-X_{l}^{(L)} X_{l}^{\dagger}\right)
\end{array}\right.
$$

whose first obvious consequence is that $\frac{\mathrm{d}}{\mathrm{d} t}\left(\hat{n}_{l}+\hat{k}_{l}\right)=0$, as we already knew from the general analysis of the integrals of motion for our model. Here we have introduced the following mean operators: $X_{l}^{(L)}=\sum_{i=1}^{L} p_{l i} X_{i}, l=1,2, \ldots, L$. Using the constant $Q_{l}=\hat{n}_{l}+\hat{k}_{l}$ and considering only the relevant equations, the above system simplifies and becomes

$$
\left\{\begin{array}{l}
\dot{X}_{l}=\mathrm{i}\left(\beta_{l}-\alpha_{l}\right) X_{l}+2 \mathrm{i} X_{l}^{(L)}\left(2 \hat{n}_{l}-Q_{l}\right)  \tag{4.5}\\
\dot{\hat{n}}_{l}=2 \mathrm{i}\left(X_{l} X_{l}^{(L)}{ }^{\dagger}-X_{l}^{(L)} X_{l}^{\dagger}\right) .
\end{array}\right.
$$

This system, as $l$ takes all the values $1,2, \ldots, L$, is a closed system of differential equations for which an unique solution surely exists. However, in order to find explicitly this solution,
it is convenient to introduce now the mean-field approximation which essentially consists in replacing the two-traders interaction $p_{i j}$ with a sort of global interaction (meaning with this that all the traders may speak among them) whose strength is inversely proportional to the number of traders: this concretely means that we have to replace $p_{i j}$ with $\frac{\tilde{p}}{L}$, with $\tilde{p} \geqslant 0$. After this replacement we have that

$$
X_{l}^{(L)}=\sum_{i=1}^{L} p_{l i} X_{i} \longrightarrow \frac{\tilde{p}}{L} \sum_{i=1}^{L} X_{i}
$$

whose limit, for $L$ diverging, only exists in suitable topologies, [6, 7], like, for instance, the strong one restricted to a set of relevant states. Let $\tau$ be such a topology. We define

$$
\begin{equation*}
X^{\infty}=\tau-\lim _{L \rightarrow \infty} \frac{\tilde{p}}{L} \sum_{i=1}^{L} X_{i} \tag{4.6}
\end{equation*}
$$

where, as it is clear, the dependence on the index $l$ is lost because of the replacement $p_{l i} \rightarrow \frac{\tilde{p}}{L}$. This is a typical behaviour of transactionally invariant quantum systems, where $p_{l, i}=p_{l-i}$. The operator $X^{\infty}$ belongs to the centre of the algebra $\mathfrak{A}$, that is it commutes with all the elements of $\mathfrak{A}:\left[X^{\infty}, A\right]=0$ for all $A \in \mathfrak{A}$. In this limit system (4.5) above becomes

$$
\left\{\begin{array}{l}
\dot{X}_{l}=\mathrm{i}\left(\beta_{l}-\alpha_{l}\right) X_{l}+2 \mathrm{i} X^{\infty}\left(2 \hat{n}_{l}-Q_{l}\right)  \tag{4.7}\\
\dot{\hat{n}}_{l}=2 \mathrm{i}\left(X_{l} X^{\infty \dagger}-X^{\infty} X_{l}^{\dagger}\right)
\end{array}\right.
$$

which, following the notation introduced in [8] in a different context, can be called the semiclassical approximation of (4.5). This system can now be solved if we assume that

$$
\begin{equation*}
\beta_{l}-\alpha_{l}=: \Phi \tag{4.8}
\end{equation*}
$$

for all $l=1,2, \ldots, L$. Under this assumption, in fact, we can deduce the time dependence of $X^{\infty}(t)$ and, as a consequence, we can completely solve system (4.7). The procedure is as follows:
(i) Using (4.7) we construct the following means: $\frac{1}{L} \sum_{l=1}^{L} \dot{X}_{l}=\frac{\mathrm{d}}{\mathrm{d} t} X_{l}^{(L)}$ and $\frac{1}{L} \sum_{l=1}^{L} \dot{\hat{n}_{l}}$.
(ii) Then we take the $\tau-\lim _{L \rightarrow \infty}$ of the system we have obtained in this way. Introducing the following intensive operators

$$
\begin{equation*}
\eta=\tau-\lim _{L \rightarrow \infty} \frac{1}{L} \sum_{l=1}^{L} \hat{n}_{l}, \quad Q=\tau-\lim _{L \rightarrow \infty} \frac{1}{L} \sum_{l=1}^{L} Q_{l}, \tag{4.9}
\end{equation*}
$$

which again belong to the centre of the algebra, we find that

$$
\left\{\begin{array}{l}
\dot{X}^{\infty}=\mathrm{i} \Phi X^{\infty}+2 \mathrm{i} X^{\infty}(2 \eta-Q)  \tag{4.10}\\
\dot{\eta}=2 \mathrm{i}\left(X^{\infty} X^{\infty \dagger}-X^{\infty} X^{\infty \dagger}\right)=0 .
\end{array}\right.
$$

This system can be easily solved: $\eta(t)=\eta$ and $X^{\infty}(t)=\mathrm{e}^{\mathrm{i} v t} X_{0}^{\infty}$, where $v=\Phi+4 \eta-2 Q$. Note that equation $\eta(t)=\eta$ has an obvious interpretation: the various $\hat{n}_{l}(t)$ change in time in such a way that their mean does not change, see (4.9). This is again a consequence of $[H, \hat{N}]=0$.
(iii) This solution must be now replaced in (4.7). It is convenient to consider two different situations: $\Phi=v$ and $\Phi \neq v$. We begin with this last case. With the change of variable $X_{l}(t)=\mathrm{e}^{\mathrm{i} t \nu}\left\{Z_{l}(t)+\frac{2}{\Phi-\nu} X_{0}^{\infty} Q_{l}\right\}$, since both $Q_{l}$ and $X_{0}^{\infty}$ do not depend on time, we deduce the following system:

$$
\left\{\begin{array}{l}
\dot{Z}_{l}=\mathrm{i}(\Phi-v) Z_{l}+4 \mathrm{i} X_{0}^{\infty} \hat{n}_{l}  \tag{4.11}\\
\dot{\hat{n}}_{l}=2 \mathrm{i}\left(Z_{l} X_{0}^{\infty \dagger}-X_{0}^{\infty} Z_{l}^{\dagger}\right),
\end{array}\right.
$$

which becomes closed if we also add the differential equation for $Z_{l}^{\dagger}$. Then we have

$$
\begin{equation*}
\dot{\Theta}_{l}(t)=\mathrm{i} \Delta \Theta_{l}(t) \tag{4.12}
\end{equation*}
$$

where we have introduced

$$
\Delta \equiv\left(\begin{array}{ccc}
\Phi-v & 4 X_{0}^{\infty} & 0 \\
2 X_{0}^{\infty} & 0 & -2 X_{0}^{\infty} \\
0 & -4 X_{0}^{\infty} & -(\Phi-v)
\end{array}\right), \quad \Theta_{l}(t) \equiv\left(\begin{array}{l}
Z_{l}(t) \\
\hat{n}_{l}(t) \\
Z_{l}^{\dagger}(t)
\end{array}\right)
$$

Remark. Note that the procedure developed here implies, as a consequence, that the dynamical behaviour of all the traders is driven by the same differential equations. This is a consequence of condition (4.8), which introduce the same quantity $\Phi$ for all the traders. Possible differences in the time evolution of the portfolios may arise therefore only because of different initial conditions. We discuss in appendix B a different approximation, which produces different equations of motion for different traders.

The solution of equation (4.12) can be written as

$$
\begin{equation*}
\Theta_{l}(t)=V \mathrm{e}^{\mathrm{i} \Delta_{d} t} V^{-1} \Theta_{l}(0) \tag{4.13}
\end{equation*}
$$

where $V$ is the matrix which diagonalizes the matrix $\Delta$ in the following sense:

$$
V^{-1} \Delta V=\Delta_{d}:=\left(\begin{array}{ccc}
\delta_{1} & 0 & 0 \\
0 & \delta_{2} & 0 \\
0 & 0 & \delta_{3}
\end{array}\right) .
$$

Remark. Note that $V$ need not to be unitary since $\Delta$ is not Hermitian.
It is clear that we are only interested in the second component of the vector $\Theta_{l}(t)$, which is exactly $\hat{n}_{l}(t)$. Carrying out all the computations and computing the mean value of $\hat{n}_{l}(t)$ on a state number $\omega_{\{n\} ;\{k\} ; O ; M}$, we find that
$n_{l}(t)=\frac{1}{\omega^{2}}\left\{n_{l}(\Phi-\nu)^{2}-8\left|X_{0}^{\infty}\right|^{2}\left(k_{l}(\cos (\omega t)-1)-n_{l}(\cos (\omega t)+1)\right)\right\}$,
where we have introduced $\omega=\sqrt{(\Phi-v)^{2}+16\left|X_{0}^{\infty}\right|^{2}}$. This formula shows that $n_{l}(t)$ is a periodic function whose period, $T=\frac{2 \pi}{\omega}$, increases when $\Phi$ approaches $v$ and when $\left|X_{0}^{\infty}\right|$ approaches zero. It is also interesting to remark that, since $\dot{n}_{l}(0)=0$ and $\ddot{n}_{l}(0)=8\left|X_{0}^{\infty}\right|^{2}\left(k_{l}-n_{l}\right), n_{l}(t)$ is an increasing function for $t$ in a right neighbourhood of 0 if $k_{l}>n_{l}$, while it is decreasing if $k_{l}<n_{l}$. This means that if $t_{l}$ has a large liquidity, then he spends money to buy shares. On the contrary, if $t_{l}$ has a lot of shares, then he tends to sell shares and to increase his liquidity, until the situation changes again.

As for the portfolio, its behaviour is the following: since $\Pi_{l}(t)=\Pi_{l}(0)+(\gamma-1)\left(n_{l}(t)-\right.$ $\left.n_{l}(0)\right)$, it is clear that $\dot{\Pi}_{l}(0)=0$ and, if $\gamma>1$ and $k_{l}>n_{l}, \ddot{\Pi}_{l}(0)>0$. This means that, in a right neighbourhood of $t=0, \Pi_{l}(t)$ increases as we expect, because of the same arguments discussed in the previous section.

Remark. It may be worth noticing that if $X_{0}^{\infty}=0$ then the number of shares does not change with time. This is a trivial consequence of (4.14) and of the definition of $\omega$, but can also be deduced directly from (4.12) and from the extremely simple expression of $\Delta$ in this case. From the first equation in (4.7) and from the definition of $\Phi$ we can also deduce that, in this case, $X_{l}(t)=\mathrm{e}^{\mathrm{i} \Phi t} X_{l}(0)$.

Let us finally consider what happens if $\Phi=v$. In this case the system (4.7) takes a simpler expression and, again, the solution can be found explicitly. Without going in many details we find $n_{l}(t)=\frac{Q_{l}}{2}+\left(n_{l}-\frac{Q_{l}}{2}\right) \cos (\omega t)+B \sin (\omega t)$, where $\omega=4\left|X_{0}^{\infty}\right|$ (since $\left.\Phi=v\right)$, and $B=\frac{2 \mathrm{i}}{\omega}\left(X_{0}^{\infty \dagger} X_{l}-X_{0}^{\infty} X_{l}^{\dagger}\right)$, which is again periodic with the same period as before.

Let us now briefly consider what happens on states of a different nature. In particular we want to understand if any meaning can be given to a KMS-state, that is, see appendix A, to an equilibrium state for a non-zero temperature.

Suppose that this is so, that is that a state $\omega_{\beta}$ satisfying condition (A.6) can be used to deduce the existence of an equilibrium for the system under consideration. It is well known that $\omega_{\beta} \neq \omega_{\{n\} ;\{k\} ; O ; M}$, so that our previous conclusions do not necessarily hold. However, if we consider the easiest nontrivial situation, $X_{0}^{\infty}=0$, it is still true that $X_{l}(t)=\mathrm{e}^{\mathrm{i} \Phi t} X_{l}=\mathrm{e}^{\mathrm{i} \Phi t} a_{l} c_{l}^{\dagger}$. If we now take $A=B^{\dagger}=X_{l}$ in (A.6), we find that $\mathrm{e}^{\beta \Phi} \omega_{\beta}\left(a_{l} a_{l}^{\dagger} c_{l}^{\dagger} c_{l}\right)=\omega_{\beta}\left(a_{l}^{\dagger} a_{l} c_{l} c_{l}^{\dagger}\right)$. Assume now that $\omega_{\beta}$ can be factorized as follows, $\omega_{\beta}=\omega_{\beta}^{(a)} \otimes \omega_{\beta}^{(c)}$, with $\omega_{\beta}^{(a)}$ related to the number of shares and $\omega_{\beta}^{(c)}$ to the cash, and let us put $m_{l}^{(a)}=\omega_{\beta}^{(a)}\left(a_{l} a_{l}^{\dagger}\right), n_{l}^{(a)}=\omega_{\beta}^{(a)}\left(a_{l}^{\dagger} a_{l}\right), m_{l}^{(c)}=\omega_{\beta}^{(c)}\left(c_{l} c_{l}^{\dagger}\right)$ and $n_{l}^{(c)}=\omega_{\beta}^{(c)}\left(c_{l}^{\dagger} c_{l}\right)$. Then the KMS condition becomes $\mathrm{e}^{\beta \Phi} m_{l}^{(a)} n_{l}^{(c)}=n_{l}^{(a)} m_{l}^{(c)}$. Since the commutation relations also imply that $m_{l}^{(a)}=1+n_{l}^{(a)}$ and $m_{l}^{(c)}=1+n_{l}^{(c)}$, this equality produces the following condition:

$$
\begin{equation*}
\mathrm{e}^{\beta \Phi}=\frac{n_{l}^{(a)}\left(1+n_{l}^{(c)}\right)}{n_{l}^{(c)}\left(1+n_{l}^{(a)}\right)}, \tag{4.15}
\end{equation*}
$$

at least if the denominator is different from zero. A first obvious remark is that, even if the single two-particles states may depend on $l$, the combination in the rhs of equation (4.15) must not.

In order to analyse condition (4.15), it is convenient to consider three different conditions: (i) $\Phi>0$, (ii) $\Phi=0$ and (iii) $\Phi<0$, and, for each of these situations, the following cases: (a) $n_{l}^{(a)}>n_{l}^{(c)}$, (b) $n_{l}^{(a)}=n_{l}^{(c)}$ or (c) $n_{l}^{(a)}<n_{l}^{(c)}$.

Case (ia). In this case, for all values of $\Phi>0$, it is not hard to check that an unique pair $\left(\beta_{0}, n_{(o)}^{(c)}\right)$ exists such that (4.15) can be verified. It is worth remarking that this also fixes the value of $n_{(o)}^{(a)}$, since $Q_{l}$ is a constant of motion. It is also possible to check that the smaller $\beta \Phi$, the larger the value of $n_{(o)}^{(c)}$, so that $n_{(o)}^{(a)}$ turns out to be smaller.

Case (ib). In this case (4.15) can be verified if and only if $\beta=0$ independently of the value of $n_{l}^{(c)}$.
Case (ic). In this case no solution of (4.15) exists.
Case (ii). In this case a solution of (4.15) exists only if $n_{l}^{(a)}=n_{l}^{(c)}$.
Finally, if $\Phi<0$, our conclusions are exactly specular to those in (i): no solution exists for (iii(a)), $\beta=0$ is the only possibility for equation (4.15) to hold in case (iii(b)) and, finally, an unique pair $\left(\beta_{0}, n_{(o)}^{(c)}\right)$ exists which verifies (4.15) in case (iii(c)).

Since $\Phi>0$ implies that $\beta_{l}>\alpha_{l}$ for all $l$, using the interpretation discussed in section 2 we could say that the inertia of the cash is larger than that of shares.

Exactly the opposite happens when $\Phi<0$, since in this case the inertia of the shares is larger than that of cash.

In $\Phi=0$ an equilibrium can be reached only if the system was already in an equilibrium state for $t=0$, i.e. if $n_{l}^{(a)}=n_{l}^{(c)}$, that is if $t_{l}$ has the same amount of cash and shares for $t=0$.

Also, if $\Phi \neq 0$ and if, for $t=0, n_{l}^{(a)}=n_{l}^{(c)}$, then an equilibrium can be reached only if $\beta=0$.

For what concerns the value of the portfolio at the time $\tilde{t}$ in which the equilibrium is reached, we get

$$
\Pi_{l}(\tilde{t})=\Pi_{l}(0)+(\gamma-1)\left(k_{l}(0)-n_{(o)}^{(c)}\right) .
$$

From this we deduce that, when $\gamma>1, t_{l}$ increments the value of his portfolio if $k_{l}(0)>n_{(o)}^{(c)}$. But, for this to be possible, the value of $\beta_{o}$ (for fixed $\Phi>0$ ) must be sufficiently high. If $\gamma<1$ the trader $t_{l}$ increments the value of his portfolio if $k_{l}(0)<n_{(o)}^{(c)}$. In this case the value of $\beta_{o}$ (again for fixed $\Phi>0$ ) must be sufficiently low.

These considerations suggest therefore to interpret $\beta^{\gamma-1}$ as a kind of information reaching the trader $t_{l}$, which should be considered together with the information already arising because of $\alpha_{l}$ and $\beta_{l}$. This is again because we are assuming that a larger amount of information produces a larger increment of the portfolio.

Remark. It must be observed, however, that in this procedure all the traders receive the same information, since $\beta^{\gamma-1}$ is the same for all $t_{l}$. What can make the difference between the traders is the information coming from $\alpha_{j}^{-1}$ and $\beta_{j}^{-1}$, as suggested in section 2. So we can distinguish between a global information, reaching all the traders in the same way, and a local information, which may be different from trader to trader.

## 5. Conclusions and outcome

In this paper we have proposed an operator approach to the analysis of some toy models of a stock market. We have shown that nontrivial results concerning the dynamical behaviour of the portfolio of each trader can be obtained, even using the existence of conserved quantities, i.e., of some operators commuting with the Hamiltonian. We have also discussed a possible use of the KMS-states within this contest.

Many things are still to be done. Among these, first of all we should introduce more than a single kind of shares. Then a different, and more realistic, mechanism to determine the price of the shares should be considered. Also, the role of condition (4.8) should be better understood, and a deeper analysis and understanding of KMS-states has to be carried out.

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## Appendix A. Mathematical background

This appendix, which is meant only for those who are not familiar with operator algebras and their applications to $Q M_{\infty}$, is essentially based on known results discussed in [9] and [10], for instance. We want to stress that only few useful facts will be discussed here, paying no particular care about the mathematical rigor. In particular we will not insist on the unbounded nature of the operators involved in the game. This is reasonable since the relevant spectrum of all the operators relevant for our discussion are usually bounded subsets of $\mathbb{R}$.

Let $\mathcal{H}$ be an Hilbert space and $B(\mathcal{H})$ the set of all the bounded operators on $\mathcal{H} . B(\mathcal{H})$ is a $\mathrm{C}^{*}$-algebra, that is an algebra with involution which is complete under a norm $\|\cdot\|$ satisfying the so-called $\mathrm{C}^{*}$-property: $\left\|A^{*} A\right\|=\|A\|^{2}$, for all $A \in B(\mathcal{H})$. As a matter of fact $B(\mathcal{H})$ is usually seen as a concrete realization of an abstract $\mathrm{C}^{*}$-algebra. It has been widely discussed
in the literature that, as far as physical applications are concerned, it is convenient to assume that the relevant observables of a certain system generate a von Neumann algebra, i.e. a closed subset of $B(\mathcal{H})$, or a topological quasi $*$-algebra. Let $\mathcal{S}$ be our physical system and $\mathfrak{A}$ the set of all the operators useful for a complete description of $\mathcal{S}$ (sometimes called the observables of $\mathcal{S}$ ). For simplicity reasons it is convenient to assume that $\mathfrak{A}$ is a $\mathrm{C}^{*}$ or a von Neumann algebra, even if this is not always possible. The description of the time evolution of $\mathcal{S}$ is related to a self-adjoint operator $H=H^{\dagger}$, which will be assumed not to depend explicitly on time, which is called the Hamiltonian of $\mathcal{S}$. Several equivalent descriptions are possible: the Schrödinger or the interaction representation, which will not be used here, or the Heisenberg representation, in which the time evolution of an observable $X \in \mathfrak{A}$ is given by

$$
\begin{equation*}
X(t)=\mathrm{e}^{\mathrm{i} H t} X \mathrm{e}^{-\mathrm{i} H t} \tag{A.1}
\end{equation*}
$$

or, equivalently, by the solution of the differential equation

$$
\begin{equation*}
\frac{\mathrm{d} X(t)}{\mathrm{d} t}=\mathrm{i} \mathrm{e}^{\mathrm{i} H t}[H, X] \mathrm{e}^{-\mathrm{i} H t}=\mathrm{i}[H, X(t)], \tag{A.2}
\end{equation*}
$$

where $[A, B]:=A B-B A$ is the commutator between $A$ and $B$. The time evolution defined in this way is usually a one-parameter group of automorphisms of $\mathfrak{A}$.

In our paper a special role is played by the so-called canonical commutation relations (CCR): we say that a set of operators $\left\{a_{l}, a_{l}^{\dagger}, l=1,2, \ldots, L\right\}$ satisfy the CCR if the following hold:

$$
\begin{equation*}
\left[a_{l}, a_{n}^{\dagger}\right]=\delta_{l n} \mathbb{I}, \quad\left[a_{l}, a_{n}\right]=\left[a_{l}^{\dagger}, a_{n}^{\dagger}\right]=0 \tag{A.3}
\end{equation*}
$$

for all $l, n=1,2, \ldots, L$. These operators, which are widely analysed in any textbook in quantum mechanics, see [11] for instance, are those which are used to describe $L$ different modes of bosons. The operators $\hat{n}_{l}=a_{l}^{\dagger} a_{l}$ and $\hat{N}=\sum_{l=1}^{L} \hat{n}_{l}$ are both self-adjoint operators. In particular $\hat{n}_{l}$ is the number operator for the lth mode, while $\hat{N}$ is the number operator of $\mathcal{S}$.

The Hilbert space of our system is constructed as follows: we introduce the vacuum of the theory, that is a vector $\varphi_{0}$ which is annihilated by all the annihilation operators $a_{l}: a_{l} \varphi_{0}=0$ for all $l=1,2, \ldots, L$. Then we act on $\varphi_{0}$ with the creation operators $a_{l}^{\dagger}$ :

$$
\begin{equation*}
\varphi_{n_{1}, n_{2}, \ldots, n_{L}}:=\frac{1}{\sqrt{n_{1}!n_{2}!\ldots n_{L}!}}\left(a_{1}^{\dagger}\right)^{n_{1}}\left(a_{2}^{\dagger}\right)^{n_{2}} \cdots\left(a_{L}^{\dagger}\right)^{n_{L}} \varphi_{0} . \tag{A.4}
\end{equation*}
$$

These vectors form an orthonormal set and are eigenstates of both $\hat{n}_{l}$ and $\hat{N}: \hat{n}_{l} \varphi_{n_{1}, n_{2}, \ldots, n_{L}}=$ $n_{l} \varphi_{n_{1}, n_{2}, \ldots, n_{L}}$ and $\hat{N} \varphi_{n_{1}, n_{2}, \ldots, n_{L}}=N \varphi_{n_{1}, n_{2}, \ldots, n_{L}}$, where $N=\sum_{l=1}^{L} n_{l}$. For this reason the following interpretation is given: if the $L$ different modes of bosons of $\mathcal{S}$ are described by the vector $\varphi_{n_{1}, n_{2}, \ldots, n_{L}}$ then $n_{1}$ bosons are in the first mode, $n_{2}$ in the second mode, and so on. The operator $\hat{n}_{l}$ acts on $\varphi_{n_{1}, n_{2}, \ldots, n_{L}}$ and returns $n_{l}$, which is exactly the number of bosons in the $l$ th mode. The operator $\hat{N}$, finally, counts the total number of bosons.

A particle in mode $l$ is created or annihilated by simply acting on $\varphi_{n_{1}, n_{2}, \ldots, n_{L}}$, respectively with $a_{l}^{\dagger}$ or $a_{l}$. Indeed we have $\hat{n}_{l}\left(a_{l} \varphi_{n_{1}, n_{2}, \ldots, n_{L}}\right)=\left(n_{l}-1\right)\left(a_{l} \varphi_{n_{1}, n_{2}, \ldots, n_{L}}\right)$ and $\hat{n}_{l}\left(a_{l}^{\dagger} \varphi_{n_{1}, n_{2}, \ldots, n_{L}}\right)=$ $\left(n_{l}+1\right)\left(a_{l}^{\dagger} \varphi_{n_{1}, n_{2}, \ldots, n_{L}}\right)$.

The Hilbert space is obtained by taking the closure of the linear span of all these vectors.
An operator $Z \in \mathfrak{A}$ is a constant of motion if it commutes with $H$. Indeed in this case equation (A.2) implies that $\dot{Z}(t)=0$, so that $Z(t)=Z$ for all $t$.

The vector $\varphi_{n_{1}, n_{2}, \ldots, n_{L}}$ in (A.4) defines a vector (or number) state over the algebra $\mathfrak{A}$ as

$$
\begin{equation*}
\omega_{n_{1}, n_{2}, \ldots, n_{L}}(X)=\left\langle\varphi_{n_{1}, n_{2}, \ldots, n_{L}}, X \varphi_{n_{1}, n_{2}, \ldots, n_{L}}\right\rangle \tag{A.5}
\end{equation*}
$$

where $\langle$,$\rangle is the scalar product in the Hilbert space \mathcal{H}$ of the theory. To be more precise, we should replace (A.5) with the following formula:

$$
\omega_{n_{1}, n_{2}, \ldots, n_{L}}(X)=\left\langle\varphi_{n_{1}, n_{2}, \ldots, n_{L}}, \pi(X) \varphi_{n_{1}, n_{2}, \ldots, n_{L}}\right\rangle
$$

where $\pi$ is a representation of the (abstract) algebra $\mathfrak{A}$ in the Hilbert space $\mathcal{H}$. We will avoid this unessential complication along this paper.

In general, a state $\omega$ over $\mathfrak{A}$ is a linear positive functional which is normalized, that is such that $\omega(\mathbb{I})=1$, where $\mathbb{I}$ is the identity of $\mathfrak{A}$. The states introduced above describe a situation in which the number of all the different modes of bosons is clear. But different states also exist and are relevant. In particular the so-called KMS state, i.e. the equilibrium state for systems with infinite degrees of freedom, are usually used to prove the existence of phase transitions or to find conditions for an equilibrium to exist. Without going into the mathematical rigorous definition, see [10], a KMS state $\omega$ with inverse temperature $\beta$ satisfies the following equality:

$$
\begin{equation*}
\omega(A B(\mathrm{i} \beta))=\omega(B A) \tag{A.6}
\end{equation*}
$$

where $A$ and $B$ are general elements of $\mathfrak{A}$ and $B(\mathrm{i} \beta)$ is the time evolution of the operator $B$ computed at the complex value $\mathrm{i} \beta$ of the time.

## Appendix B. On system (4.7)

We will show now how to solve system (4.7) without using condition (4.8).
For this we introduce the following quantities: $\gamma_{l}=\beta_{l}-\alpha_{l}, X_{\gamma^{k}}^{\infty}=\tau-\lim _{L} \frac{1}{L} \sum_{l=1}^{L} \gamma_{l}^{k}$ $X_{l}, k=1,2, \ldots, \eta_{\gamma}=\tau-\lim _{L} \frac{1}{L} \sum_{l=1}^{L} \gamma_{l} \hat{n}_{l}$, and $Q_{\gamma}=\tau-\lim _{L} \frac{1}{L} \sum_{l=1}^{L} \gamma_{l} Q_{l}$. Of course, we are assuming here that all these limits do exist. Repeating the same steps as in section 4 , we find the following system:

$$
\left\{\begin{array}{l}
\dot{X}^{\infty}=\mathrm{i} X_{\gamma}^{\infty}+2 \mathrm{i} X^{\infty}(2 \eta-Q) \\
\dot{\eta}=0
\end{array}\right.
$$

To close this system, we also need the differential equation for $X_{\gamma}^{\infty}$ which, as it is easily understood, involves $X_{\gamma^{2}}^{\infty}, \eta_{\gamma}$ and $Q_{\gamma}$. Note that, in our previous approximation, these operators turned out to be equal respectively to $\Phi^{2} X^{\infty}, \Phi \eta$ and $\Phi Q$. Moreover, in that approximation, we also had $X_{\gamma}^{\infty}=\Phi X^{\infty}$, so that the system above was already closed. We improve our original approximation by taking now $X_{\gamma}^{\infty}$ as a new variable and replacing $X_{\gamma^{2}}^{\infty}, \eta_{\gamma}$ and $Q_{\gamma}$ with $\tilde{\Phi}^{2} X^{\infty}, \tilde{\Phi} \eta$ and $\tilde{\Phi} Q$, where we have introduced $\tilde{\Phi}=\lim _{L} \frac{1}{L} \sum_{l=1}^{L} \gamma_{l}$, assuming that it exists. It should be noticed that $\tilde{\Phi}$ extends $\Phi$ in the sense that they coincide if $\gamma_{l}=\Phi$ for all $l$. The equation for $X_{\gamma}^{\infty}$ is therefore $\dot{X}_{\gamma}^{\infty}=\mathrm{i} \tilde{\Phi}^{2} X^{\infty}+2 \mathrm{i} X^{\infty} \tilde{\Phi}(2 \eta-Q)$. For the sake of simplicity we will work here assuming that $X_{\gamma}^{\infty}(0)=0$ and $2 \mu+\tilde{\Phi}=0$, where $\mu=2 \eta-Q$. With these assumptions, which could be avoided in a more general analysis, we can repeat the same steps as in section 4 , getting the following result:
$n_{l}(t)=\frac{1}{\omega_{l}^{2}}\left\{n_{l}\left(\gamma_{l}+\tilde{\Phi}\right)^{2}-\frac{32 \mu^{2}}{\tilde{\Phi}^{2}}\left|X_{0}^{\infty}\right|^{2}\left(k_{l}\left(\cos \left(\omega_{l} t\right)-1\right)-n_{l}\left(\cos \left(\omega_{l} t\right)+1\right)\right)\right\}$,
where we have introduced $\omega_{l}=\sqrt{\left(\gamma_{l}+\tilde{\Phi}\right)^{2}+\frac{64 \mu^{2}}{\tilde{\Phi}^{2}}\left|X_{0}^{\infty}\right|^{2}}$. It is clear now that different traders may have different behaviours, depending on the related value of $\gamma_{l}$ : it is interesting to notice, for instance, that if $\left|\gamma_{l}\right| \rightarrow \infty$, that is when $\alpha_{l}$ and $\beta_{l}$ are very different from each other, then $n_{l}(t)=n_{l}$. This is not so for zero or intermediate values of $\left|\gamma_{l}\right|$, for which a nontrivial time evolution of $n_{l}(t)$ is recovered.

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